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Magnetic and petrologic characterization of synthetic Martian basalts and implications for the surface magnetization of Mars

Julie A. Bowles,1,2 Julia E. Hammer,1 and Stefanie A. Brachfeld3

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[1] A suite of synthetic Martian basalts is generated with the objective of providing fundamental material properties data for use in modeling and interpretation of mission data. We systematically evaluate the effects of major element composition, oxygen fugacity (fO2), and cooling rate on phase chemistry and magnetic mineralogy, grain size, and intensity of remanent magnetization. The range of experimental compositions and fO2 are chosen to bracket the range expected in the Martian crust; our results should therefore span the range of possible mineralogies, textures, and magnetic properties in rapidly cooled Mars crustal materials. Two starting compositions are used for the sample synthesis: (1) an Fe-rich, Al-poor composition patterned after SNC basaltic meteorites and (2) a composition based on thermal emission spectrometer (TES) data with a much lower Fe/Al ratio. The resulting magnetic phase in samples generated at the quartz-fayalite-magnetite (QFM) buffer is a spinel-structured oxide with varying amounts of Cr, Ti, Mg, and Al. Compositional differences depend on bulk composition, cooling rate, differences in crystallization sequence, and the kinetics of silicate mineral nucleation and growth. Oxide abundance and magnetic intensity are most strongly influenced by fO2, with more subtle composition and cooling rate effects. Moderately oxidizing QFM conditions result in an intense magnetization (2.3 × 10⁻⁵ Am² kg⁻¹ to 1.4 × 10⁻² Am² kg⁻¹), especially in the meteorite-derived basalts. However, an increase of magnetic grain size into the multidomain range (meteorite-type) and/or low unblocking temperatures resulting from increased Cr substitution (TES-type) may affect the long-term stability of the remanence in QFM samples.


1. Introduction

[2] The strong intensity of the Martian magnetic anomalies mapped by the Mars Global Surveyor (MGS) has led to considerable interest in the magnetization of the Martian crust. The strongly magnetic crust is limited largely to the ancient (>4 Ga), cratered, southern highlands, and the standard interpretation holds that the crust was magnetized in the presence of a Martian dynamo that was active early in the planet’s history [e.g., Stevenson, 2001]. A growing body of work seeks to further explain the patterns of magnetization [Johnson and Phillips, 2005; Nimmo, 2000; Stanley et al., 2008; Hood et al., 2007], as well as a magnetic mineral assemblage capable of producing a magnetization inferred to be several orders of magnitude stronger than typical terrestrial basalt [Connerney et al., 1999].

[3] Most rock magnetic studies have focused on the magnetic mineralogy of the geologically youthful (<1.2 Ga) SNC meteorites, which postdate cessation of the Martian geodynamo and are inferred to be from the northern lowlands. Magnetic remanence carriers in these meteorites include pyrrhotite, (titanom)agnetite, and chromite [Rochette et al., 2005, and references therein; Yu and Gee, 2005]. At least two meteorites have been found to contain metallic iron inclusions within olivine crystals, an inferred shock-induced reduction product [Van de Moortele et al., 2007]. Of the SNC meteorites, typically only the nakhlites (clinopyroxenites) and some basaltic shergottites are capable of acquiring a remanence intense enough to produce anomalies of sufficient magnitude [Rochette et al., 2005].

[4] Because the magnetic anomalies are largely limited to the more ancient southern highlands, these young meteorites may be atypical of the magnetized crust. Mineral phases not observed in meteorites have therefore also been proposed, including multidomain hematite [Kletetschka et al., 2000; Dunlop and Kletetschka, 2001]; olivine-hosted magnetite inclusions [Gunlaugsson et al., 2006]; and hematite-ilmenite intergrowths [McEnroe et al., 2002; Brown and McEnroe, 2008] or some other two-phase intergrowth [Hargraves et al., 2001].
[5] While both shock [Van de Moortêle et al., 2007] and hydrothermal alteration [Scott and Fuller, 2004] have been suggested to explain the intense magnetization, most theories posit that the anomalies represent a primary thermal remanence acquired by igneous rocks cooling in presence of a dynamo field. This in part has motivated us to explore what conditions might be necessary to produce a strong, stable magnetic carrier in a solidifying Martian crust. We take an experimental approach, synthesizing basalts of expected Martian compositions, while varying cooling rate, oxygen fugacity ($f_O^2$), and major element composition. This allows us to place constraints on conditions that could plausibly result in strong, stable magnetic remanence, and provides fundamental material properties for use in modeling and interpretation of mission data.

[6] Our work builds upon that of Hammer [2006], which suggested that the anomalies represent a primary thermal remanence acquired by igneous rocks cooling in presence of a dynamo field. This in part has motivated us to explore what conditions might be necessary to produce a strong, stable magnetic carrier in a solidifying Martian crust. We take an experimental approach, synthesizing basalts of expected Martian compositions, while varying cooling rate, oxygen fugacity ($f_O^2$), and major element composition. This allows us to place constraints on conditions that could plausibly result in strong, stable magnetic remanence, and provides fundamental material properties for use in modeling and interpretation of mission data.

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Table 1. Bulk Compositions and CIPW Normative Mineralogy

<table>
<thead>
<tr>
<th></th>
<th>Meteorite Type</th>
<th>TES Type</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>A* $^{a}$</td>
<td>A* mod$^{b}$</td>
</tr>
<tr>
<td>SiO$_2$</td>
<td>51.40</td>
<td>51.16</td>
</tr>
<tr>
<td>TiO$_2$</td>
<td>1.63</td>
<td>1.61</td>
</tr>
<tr>
<td>Al$_2$O$_3$</td>
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<td>8.70</td>
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<tr>
<td>FeO*</td>
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<td>18.90</td>
</tr>
<tr>
<td>MgO</td>
<td>7.27</td>
<td>7.04</td>
</tr>
<tr>
<td>CaO</td>
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</tr>
<tr>
<td>Na$_2$O</td>
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<td>2.28</td>
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<tr>
<td>K$_2$O</td>
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<td>0.76</td>
</tr>
<tr>
<td>P$_2$O$_5$</td>
<td>0.49</td>
<td>0.50</td>
</tr>
<tr>
<td>MnO</td>
<td>–</td>
<td>0.53</td>
</tr>
<tr>
<td>Cr$_2$O$_3$</td>
<td>0.10</td>
<td>0.10</td>
</tr>
<tr>
<td>Fe/Al</td>
<td>1.50</td>
<td>1.50</td>
</tr>
</tbody>
</table>

Normative mineralogy$^{h}$

- Quartz 4.84 4.10 7.65 4.88
- Total feldspar/feldspathoid 35.25 34.86 34.39 32.89
- Total pyroxene 49.26 49.78 47.53 51.48
- Olivine 0.00 0.00 0.00 0.00
- Apatite 1.15 1.18 1.08 0.78
- Ilmenite 3.07 3.04 2.84 3.04
- Magnetite 6.80 7.05 6.51 6.93

$^{a}$SNC-like composition used by Hammer [2006] and Brachfeld and Hammer [2006].
$^{b}$Desired SNC-like composition used in present study which includes Mn and Cr.
$^{c}$Measured bulk composition of glass quenched from >T$_L$.
$^{d}$Desired Earth-like composition used in present study.
$^{e}$Measured bulk composition of glass quenched from >T$_L$. Note that the SiO$_2$ content of Batch 2 was considerably lower than desired, resulting in a melt that is olivine normative rather than quartz normative. It is possible this had an effect on the silicate phases present, although it is not obvious in the experimental results. It is unlikely to play a significant role in the abundance or composition of oxide phases.
$^{f}$McSween et al. [2006], as reproduced in the work by Nekvasil et al. [2009].
$^{g}$Did not analyze for Cr.
$^{h}$FeO and Fe$_2$O$_3$ were computed from Fe* using method of Kilinc et al. [1983] and assuming 1200°C and fO$_2$ controlled by the QFM buffer in order to perform the CIPW norm calculation. Normative mineralogy for Backstay from McSween et al. [2006].

2. Methods

2.1. Starting Material Selection

[8] Two starting compositions (Table 1) are used for the sample synthesis. The first composition is Fe-rich and Al-poor (atoms Fe/Al = 1.5) and is patterned after the widely recognized Fe-rich, Al-poor character of SNC parent melts [e.g., Johnson et al., 1991; Treiman, 1993; McSween, 2007]. It is similar in composition to that used by Hammer [2006] and Brachfeld and Hammer [2006], but includes small amounts of Mn and Cr. The second composition is identical to the first in all respects save Fe (in which the second composition is poorer), Al (richer), and Ca (slightly richer). The relative abundances of Fe and Al in the second composition (Fe/Al = 0.4) typify those of terrestrial basalt.

iron-wüstite (IW) buffer and the quartz-fayalite-magnetite (QFM) buffer. This allows us to explore in more detail the $f_O^2$ range most commonly found for basaltic meteorites [Her et al., 2002; Wadhwa, 2008], as well as the range over which a dramatic change in both oxide abundance and magnetic properties occurs [Brachfeld and Hammer, 2006].

[7] While both shock [Van de Moortêle et al., 2007] and hydrothermal alteration [Scott and Fuller, 2004] have been suggested to explain the intense magnetization, most theories posit that the anomalies represent a primary thermal remanence acquired by igneous rocks cooling in presence of a dynamo field. This in part has motivated us to explore what conditions might be necessary to produce a strong, stable magnetic carrier in a solidifying Martian crust. We take an experimental approach, synthesizing basalts of expected Martian compositions, while varying cooling rate, oxygen fugacity ($f_O^2$), and major element composition. This allows us to place constraints on conditions that could plausibly result in strong, stable magnetic remanence, and provides fundamental material properties for use in modeling and interpretation of mission data.

[6] Our work builds upon that of Hammer [2006] and Brachfeld and Hammer [2006] where a synthetic, Fe-rich basalt (patterned after the basaltic meteorites) was crystallized over a wide range of $f_O^2$ and cooling rates, and the resulting magnetic assemblage was examined. In this work, we expand the range of compositions to include a basalt with a much lower Fe/Al ratio, on the basis of satellite thermal emission spectrometer (TES) observations of the southern highlands that suggest a composition more closely resembling terrestrial basalts [Hamilton et al., 2001]. We also add the minor elements Mn and Cr to both the Fe-rich and Fe-poor basalt, in light of evidence that Cr-Fe-Ti oxides can carry a magnetic remanence in Martian meteorites [Yu and Gee, 2005; Weiss et al., 2002]. Compared to Hammer [2006], we restrict the range of $f_O^2$ to that between the
[e.g., Leeman et al., 1990; Flower et al., 1982] and satisfy compositional constraints posed by deconvolved thermal emission spectrometer (TES) data from the Mars surface [e.g., Hamilton et al., 2001]. Some experimentalists have also proposed that moderate-Fe/Al basalts are parental to Mars meteorite materials [Nekvasil et al., 2007]. Thus, the experimental starting compositions synthesized here bracket the range of compositions that have been proposed as relevant mafic igneous materials on Mars.

[9] The first composition (Fe/Al = 1.5) will hereafter be referred to as the “meteorite-type” or “M-type” composition, while the second (Fe/Al = 0.4) will be referred to as the “TES-type” or “T-type” composition.

[10] The CIPW normative mineralogy (Table 1) and rheologies of the M-type and T-type target starting materials differ in important ways. Key differences include (1) the inverse relationship between feldspar and pyroxene contents, with 35 and 50 wt. %, respectively, in the M-type composition, and ~60 and 35 wt. %, respectively, in the T-type composition; (2) the normative magnetite content of the M-type material more than twice that of T-type material; and (3) a melt viscosity more than 1 order of magnitude greater in T-type melt compared to M-type melt (using method of Shaw [1972]), at a given temperature.

[11] Differences in the major element composition of the two starting materials lead to contrasting liquid lines of descent during equilibrium fractional crystallization as computed using the MELTS algorithm [Ghiorso and Sack, 1995]. Consistent with the CIPW norm calculation, the mineral-in-sequence at QFM in the M-type liquid is spinel oxide, clinopyroxene, feldspar, and olivine; the T-type liquid begins with spinel oxide, followed by feldspar, olivine, and clinopyroxene. Similar trends are computed for reducing conditions, with a key difference being the increase in the olivine-in temperature (becoming the liquidus phase) and greater modal proportion of olivine in M-type material with decreasing fO2.

2.2. Sample Synthesis

[12] Samples were synthesized from powdered oxide, carbonate, and phosphate reagents. This was done twice for each composition, as we ran out of the first batch midway through the experiments. Slight differences in the resulting bulk composition are noted in Table 1. Oxygen fugacity was maintained in one of the following two ways.

[13] 1. Sample powder mixtures intended for treatment at QFM and some intended for IW were first heated at 925°C in a controlled fO2 atmosphere using flowing H2-CO2 to impart an intrinsic fO2. Approximately 0.3 g of the mixture was then placed in a Pt capsule, which was sealed in an evacuated Si tube along with a solid state buffer assemblage to monitor departure from the desired fO2. QFM samples at cooling rates ≤ 19°C h–1 were sealed in double-walled Si tubes, on the basis of results from Hammer [2006] showing unintended oxidation to the magnetite-hematite buffer.

[14] 2. Some samples to be treated at IW and all at intermediate fO2 were heated to 925°C in an arbitrary atmosphere to decarbonate the sample powder mixtures. About 0.3 g of the decarbonated sample mixture was then pressed into a pellet and was suspended on a Pt wire spiral to be heated above the liquidus (Tli) under flowing H2-CO2.

[15] In both scenarios, fO2 during flowing gas treatment was monitored with a solid zirconia electrolyte oxygen sensor. Temperature directly adjacent to the sample was measured with a Pt-Rh thermocouple.

[16] To address the issue of iron loss into the Pt capsule or Pt wire, capsules to be used at fO2 ≥ 1 × 1017 to 2 × 1017 units were run first with the sample material at Tli + 5°C for 18 h and then quenched. The solidified sample was discarded and this procedure was repeated one or more times [e.g., Hammer, 2006]. Because Fe loss increases at decreasing fO2, a more aggressive approach was used for samples run at IW. The equilibrium mole fraction of Fe in the Pt (XFe) was calculated using the equations of Kessel et al. [2001] for temperatures just below the liquidus. The Pt was then packed in Fe2O3 powder and held at the temperature and fO2 determined to produce the desired XFe [Kessel et al., 2001]. Because wüstite (FeO) adhered to the Pt following this treatment and had to be abraded away, it was impossible to determine the precise value of XFe via weighing, but we estimate that XFe is within ~20% of the desired value.

[17] Most samples were heated to ~5–10°C above Tli as calculated the MELTS algorithm [Ghiorso and Sack, 1995] and were held at this temperature for 4–6 h. Alternatively, a few samples were heated to ~Tli + 25°C and were cooled at 2°C h–1 to Tli + 5°C where they were held for 4 h. Samples were then cooled to 400–600°C at constant cooling rates ranging from 231°C h–1 to 5.7°C h–1 using a programmable temperature controller. In addition, an exponentially varying cooling rate was also applied to QFM samples to simulate the natural cooling path of a conductively cooling lava flow at ~0.25 m from its surface [Turcotte and Schubert, 2002]. Cooling rates during this experiment varied from 53°C h–1 at high temperatures to 2°C h–1 at the lowest temperatures.

[18] Following sample synthesis, splits of the recovered material were used for compositional, petrographic and magnetic analyses.

2.3. Compositional and Petrographic Analyses

[19] Electron backscatter images and nonquantitative EDS constraints on mineral compositions were acquired with a JEOL JSM900 scanning electron microscope (SEM) at University of Hawaii. Compositional information on the largest oxide crystals was obtained using a Cameca SX-100 electron microprobe (EMP) at University of Oregon.

[20] Low totals were consistently observed on the EMP analyses, and it was ultimately determined that this resulted both from a carbon coating thinner than that used on the calibration standards and (to a lesser extent and in the smallest grains only) an interaction volume larger than the grain size. Some samples were measured again after reapplying an appropriate carbon coat. After recoating, analysis totals were near 100%, assuming 4 oxygen per 3 cations (spinel structure). If analyses are all normalized to 100%, the within sample variation is considerably larger than variation because of either of these two effects; we therefore consider all analyses in determining relative compositional variability between samples. Contamination of analyses by the surrounding matrix was monitored by measuring Si content, and any analysis with >0.5 wt% SiO2 was discarded.
[21] To characterize oxide compositional variation among samples, we assumed spinel structure for all analyses, and calculated spinel end-members in the following manner:

[22] 1. All Ti was assigned to ulvöspinel (Fe$_2$TiO$_4$).
[23] 2. All Mn (typically <0.5 wt% MnO) was assigned to jacobite (MnFe$_2$O$_4$).
[24] 3. Al and Mg are assigned first to spinel (MgAl$_2$O$_4$), then excess Al was assigned to hercynite (FeAl$_2$O$_4$) or excess Mg was assigned to magnesiowustite (Mg$_2$FeO$_4$), then to magnesiochromite (MgCr$_2$O$_4$) if all Fe was depleted.
[25] 4. The remaining Cr was assigned to chromite (FeCr$_2$O$_4$).
[26] 5. The remaining Fe was assigned to magnetite (Fe$_3$O$_4$).

[27] While this is a somewhat complex and nonunique solution, fewer end-members could not accommodate the often large amounts of Mg, Al, and Cr.

[28] The area fractions of select phases were measured by thresholding the characteristic range of gray-scale values in backscatter electron images using ImageJ freeware (http://rsb.info.nih.gov/ij/). For materials lacking textural fabric, the volume fraction of a given phase is equivalent to its area fraction in a 2-D slice [Hilliard, 1968].

### 2.4. Rock Magnetic Analyses

[29] Rock magnetic measurements were designed to place constraints on the grain size and composition of the magnetic phases present, as well as to test the remanence-bearing capacity of the samples. Two sample splits were used for magnetic analyses. The first was used to measure saturation magnetization ($M_s$) as a function of temperature from 50°C to 600°C in flowing He gas in an attempt to estimate Curie temperature ($T_c$). A more complete battery of tests was carried out on the second split. Acquisition of anhysteretic remanent magnetization (ARM) was measured for all samples in a 100 mT peak alternating field (AF) and a maximum 200 μT DC bias field; this was followed by AF demagnetization of the ARM. Magnetization (M) versus applied field (H) data were acquired on a subset of samples in order to determine hysteresis behavior of the samples for applied field (H) data were acquired on a subset of samples in order to determine hysteresis behavior of the samples for temperatures between 10 K and 300 K. Field cooled – zero field cooled (FC-ZFC) experiments were carried out in part to constrain magnetic mineralogy: a 2.5 T saturation isothermal remanent magnetization (SIRM) was applied at 20 K and then measured while warming to 300 K following cooling in zero field and in a 2.5 T field. Additionally, on some samples a room temperature (300 K) SIRM was measured on cooling to 20 K and on warming back to room temperature. Susceptibility as a function of temperature (20–300 K) and frequency was also measured on many samples. Finally, samples were given a thermal remanence (TRM) at 550°C (chosen on the basis of $T_c$) in a 49 μT field in Ar atmosphere. The TRM was measured and then progressively removed in 50°C steps between 100 and 550°C. To monitor alteration, low-field susceptibility was measured following each heating treatment, and a second TRM at 49 μT was given following the complete thermal demagnetization. Finally, TRMs were additionally imparted in 35 μT and 80 μT fields. All room temperature remanence measurements (ARM, TRM) were made in the paleomagnetic laboratory at University of Hawai‘i on a 2 G cryogenic rock magnetometer. The remainder of the measurements was made at the Institute for Rock Magnetism at the University of Minnesota.

### 3. Results

[30] Where used, the buffer assemblage indicators showed no evidence for significant departure from either the IW or QFM buffer. In the more slowly cooled QFM samples, the assemblage was dominated by quartz and magnetite; although small amounts of fayalite could be found, it was typically armored by quartz grains, suggesting that the oxygen fugacity may have become somewhat more oxidizing than QFM. However, one experiment (T-type, 5.7°C h$^{-1}$) was repeated under flowing gases while monitoring $f$O$_2$, and oxide compositions as measured by EMP were entirely consistent with those in the sample cooled in vacuum with the buffer assemblage.

#### 3.1. Petrographic and Compositional Analyses

[31] Many of the charges are texturally heterogeneous on the millimeter scale, with spatial gradients in crystallinity, crystal size and morphology, and to a lesser extent, mineralogy. These variations are consistent with heterogeneous nucleation and diffusion controlled growth occurring during rapid cooling. Portions of material from throughout the charge were included in all thin sections and textural characterizations.

[32] The dominant phase in all experiments, regardless of bulk composition, cooling rate or $f$O$_2$, is Ca-rich pyroxene (Figures 1 and 2). It occurs in a spectacular range of morphologies, ranging from feathery dendrites (e.g., Figures 1f, 1g, and 2g) to zoned, faceted laths (Figures 2c and 2d); in T-type samples it is frequently intimately intergrown with plagioclase (Figures 2f and 2h), a phase not observed in the M-type samples. Olivine (Figures 1a, 1c, and 1d) is present in the M-type samples (most prevalent at IW), while only two T-type samples (IW; 18.7 and 5.7°C h$^{-1}$) have olivine (not shown). The T-type samples were on average finer grained and more crystalline than the M-type samples. Observations and trends in mineral abundance and texture with respect to $f$O$_2$ and cooling rate are detailed below.

##### 3.1.1. Effects of Cooling Rate and $f$O$_2$ on Silicate Minerals

[33] Predictably, the result of decreasing cooling rate on both M- and T-type samples is to coarsen and mature the morphologies of silicate (and oxide) mineral phases. At 231°C h$^{-1}$, the morphologies of pyroxene and olivine formed at all $f$O$_2$ values are dendritic, with large specific surface areas and low number densities suggesting that the dominant crystallization mechanism is diffusion-limited crystal growth. Plagioclase and Ca pyroxene in the holo-crystalline T-type samples formed during rapid cooling (231°C h$^{-1}$) resemble fine-scale eutectic intergrowths (similar to texture shown in Figure 1e). Compact and faceted crystals, which presumably formed by interface-limited crystal growth [Kirkpatrick, 1981], occur only in slow-cooled experiments (e.g., compare Figures 2a and 2d).

[34] Of the two samples that followed an exponentially varying temperature-time path, the M-type sample has pyroxene morphologies most similar to those of the more rapidly cooled (≥72.4°C h$^{-1}$) samples (Figure 2e). The T-type sample is extremely heterogeneous; portions of the
sample are wholly glass, while other portions contain feathery or dendritic pyroxene, and some areas contain more faceted pyroxene (Figure 2j).

[35] The effects of varying $fO_2$ on M-type silicates are broadly similar at all cooling rates, chiefly controlling olivine abundance and morphology. The olivine correlation is weak at the fastest cooling rate (231°C h$^{-1}$), while at 72°C h$^{-1}$, an increase in $fO_2$ from IW to QFM is accompanied by maturation of olivine morphology from skeletal toward faceted forms at essentially constant volumetric abundance. At slower cooling rates (18.7 and 5.7°C h$^{-1}$), increasing $fO_2$ is associated with decreased olivine content (from 15 and 8 vol.%, respectively, at IW, to 0 vol.% at QFM).

[36] In contrast, the effect of increasing $fO_2$ is incoherent with respect to silicate mineralogy and morphology in T-type experiments. With slow cooling (18.7 and 5.7°C h$^{-1}$), increasing $fO_2$ is associated with declining olivine abundance (as observed in the M-type runs) as well as dramatically increasing plagioclase abundance (from 0 to >30 vol.%). However, with fast cooling (231 and 72.4°C h$^{-1}$), olivine content is unchanged and plagioclase is either unchanged (at ~45 vol.%) or completely absent (i.e., 0 vol.%) with increasing $fO_2$.

3.1.2. Effects of Cooling Rate and $fO_2$ on Oxide Mineralogy

[37] Similar to the silicate phases, the effect of decreasing cooling rate at QFM in M-type oxides is to evolve from cruciform and skeletal oxides (Figure 2a, 2b, and 2c) to more equant and faceted grains (Figure 2d). Additionally, oxide grain size increases, from maximum grain diameters of ~8 μm at 231°C h$^{-1}$ to ~40 μm at 5.7°C h$^{-1}$. The variable cooling rate sample has a set of faceted grains with maximum diameters approaching that of the most slowly

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Figure 1. (a–g) Electron backscatter images of samples cooled at 72 K h$^{-1}$ at indicated $fO_2$. There is a general (but not absolute) trend toward larger and more abundant oxides at higher $fO_2$. Here cpx, clinopyroxene; ol, olivine; plag, plagioclase; tmt, titanomagnetite; Fe(-Ti)-Cr ox, iron(-titanium)-chromium oxide. All oxides also contain some Mg + Al. Scale bar shown in Figure 1g applies to all other images.
cooled sample. However, this sample also has a population of smaller, cruciform grains that are not present in the most slowly cooled sample; most of these cruciform oxides are observed decorating pyroxene grain boundaries (Figure 2e).

[38] Increasing $\phi$O$_2$ in M-type samples from IW to QFM at all cooling rates causes oxide abundance to increase. At $231^\circ$C h$^{-1}$, the increase is 0 to 4 vol.%; at $72^\circ$C h$^{-1}$, 0 to 3 vol.%; at $18.7^\circ$C h$^{-1}$, 1 to 4 vol.%; and at $5.7^\circ$C h$^{-1}$, 0 to 8 vol.%. In the more reduced samples, submicron oxides (when present) frequently appear either as decorations or as irregular rims on edges of pyroxene crystals. Two of the intermediate $\phi$O$_2$ samples (IW+2, $18.7^\circ$C h$^{-1}$; IW + 2.75, $72.4^\circ$C h$^{-1}$) contain equant (~1–2 $\mu$m diameter), Cr-rich oxides within olivine hoppers (Figure 1c). Many samples

Figure 2. (a–j) Electron backscatter images of samples cooled under QFM conditions at indicated cooling rates. Here “var” indicates exponentially varying cooling path. Scale bar is 50 $\mu$m. There is a general (but not absolute) trend toward larger grain size at slower cooling rates, most prominent in the M-type samples. Oxides appear as the highly reflective (bright) phase in each image. Figure 2b also contains numerous skeletal and cruciform oxides. The fine-grained groundmass of Figures 2f, 2h, and 2i appears to be an intergrowth of plagioclase and clinopyroxene with interstitial oxides. Here gl, glass. Scale bar shown in Figure 2f applies to all other images.
and cooling rates, Mt is uMt e at ftp://ftp.agu.org/apend/je/xx below the QFM buffer, as well as smaller grains in m BOWLES ET AL.: MAGNETIC PROPERTIES OF MARTIAN BASALTS (where Mt Mt is mole fraction of Z m Mt is.)

Results of the spinel and 2d) in this composition.

may contain lower-than-expected oxide contents because of greater capacity for oxide formation, the M-type material in the cooling history. Despite its greater Fe content and thus double the normative magnetite content (Table 1), and reflected in Figure 1), despite the M-type material having a very low degree of crystallinity and maximum grain diameters of ~0.5 μm. The clusters of Cr-rich oxides observed in M-type samples are also present in T-type samples, again more frequently under oxidizing conditions.

3.1.3. Comparison of M-Type and T-Type Mineralogy

[40] One important difference between the M-type and T-type experiments is the degree of crystallinity, which is consistently lower in the M-type samples. The largest difference is between experiments performed during slow cooling at QFM: the T-type sample is holocrystalline (100 vol.% crystals (Figure 2i)), while the M-type sample contains only ~50 vol.% crystals (Figure 2d). Moreover, none of the M-type samples is more than 60 vol.% crystalline. The difference is smaller (but not consistently so) at reducing conditions and more rapid cooling rates. This difference in crystallinity relates to two other observations: (1) oxides are equally abundant (~8 vol.% in slow-cooled M-type and T-type runs at QFM (although this is not reflected in Figure 1), despite the M-type material having double the normative magnetite content (Table 1), and (2) samples that reached high crystallinities (>80 vol.% are invariably plagioclase-rich. Pertinent to both points, textures in the holocrystalline T-type runs (Figures 2h and 2i) suggest that extensive silicate crystallization in these experiments caused interstitial oxide crystallization relatively late in the cooling history. Despite its greater Fe content and thus greater capacity for oxide formation, the M-type material may contain lower-than-expected oxide contents because of comparatively sluggish silicate crystallization (Figures 2c and 2d) in this composition.

3.1.4. Oxide Compositions

[41] Quantitative compositional information by electron microprobe analysis (EMP) was only possible on grains >1 μm. This precluded analyses on most samples formed at fO2 below the QFM buffer, as well as smaller grains in QFM samples. These smallest grains probably differ in composition from the larger crystals, insofar as size is correlated with the timing of nucleation with respect to cooling and therefore melt composition. With these caveats in mind, we report (Data Set S1) EMP analyses for all of the samples cooled at QFM (with the exception of the variable cooling rate T-type sample).1 Results of the spinel end-member calculations are shown in Figure 3.

[42] To a large degree, the compositional trends in the experimentally grown oxides are consistent with expectation based on compositional differences between the M-type and T-type starting materials. For example, the oxides in Al-poor M-type run products have considerably less Al than the oxides in Al-rich T-type material. For crystals with nominal Cr content (~0.02 atoms per formula unit), the M-type material crystallizes oxides with Al content of ~0.14 atoms per formula unit (apfu) compared to ~0.45 for the T-type starting material. Both starting materials produce oxides with a wide variety of chromite contents (0 to 1.0 apfu Cr), tending to decrease as crystallization progresses, as ascertained from zoned crystals (Cr content decreases from core to rim). Such Cr zoning in spinels is frequently observed in nature [e.g., Haggerty, 1991; de Wall et al., 2004].

[43] The slowest-cooled and variable cooling rate M-type samples are the only charges to exhibit compositionally homogeneous oxides. These faceted, coarse-grained crystals cluster near Mt53Usp12Z10 (where Mt49 is mole fraction of magnetite, Usp12 is ulvospinel, and Z10 represents the summed mole fractions of spinel, chromite, magnesiowüstite, magnesiocromite, and jaspilite compositional end-members) regardless of the size of the crystal analyzed, or placement of the focused electron beam with respect to the crystal core (Figure 3a). These crystals have roughly equal amounts of Ti, Mg, and Al and are relatively low in Cr (Figure 3b).

[44] M-type samples cooled at moderate rates exhibit a population composition near the slow-cooled cluster (with a slightly higher ulvospinel content) as well as sparse crystals that are much richer in non titanomagnetcite components (ranging up to Mt50Usp30Z10). The molar proportions of Mg end-members are reduced in most analyses compared to those in the slow-cooled samples, and much of the scatter in the data can be explained by varying Cr concentration. The most rapidly cooled M-type oxides span a continuous range between Mt53Usp12Z10 and Mt30Usp10Z70, with no significant gaps in the spectrum. The spread in the data is almost entirely explained by variations in the Cr content (from 0.00 to 0.45 mol fraction chromite), and to a lesser extent, Mg and Al end-member components.

[45] In contrast to the M-type experiments, the slowest cooled T-type run at QFM contains a wide variety of oxide compositions. Resembling the fastest-cooled M-type run, the crystals analyzed in this sample exhibit a broad, continuous array of non titanomagnetcite contents from Mt50Usp12Z15 and Mt20Usp70Z75. Similar to the M-type run, most of the spread in the data can be explained by magnetite and ulvospinel components being replaced by a chromite component, and to a lesser extent by incorporation of Mg and Al end-member components. The samples cooled at the intermediate and fastest rates contain compositionally discrete crystals (rather than a continuum), with grains clustering near Mt7Usp2Z80 and Mt3Usp5Z90. These compositional differences have no obvious correlation with crystal size or morphology.

[46] Scarcity of crystals coarse enough for quantitative analysis suggests that the compositions plotted in Figure 3 may not accurately reflect the complete range and distribution of compositions in these samples. However, uniformity of the compositional trends (including the common Mt-Usp

ratio among Usp-poor crystals and clustering of Usp-rich compositions) suggests that a similar liquid line of descent was traversed by all samples in the set, albeit to varying extents depending upon cooling rate.

3.2. Magnetic Results

3.2.1. M_s Versus T (Curie Temperature Estimates)

Curie temperature ($T_c$) was determined from the $M_s$ versus $T$ measurements (by the method of two intersecting tangents) on all meteorite-type QFM samples, and two of five TES-type QFM samples. In other samples, the paramagnetic signal dominated the measurements. In the M-type samples, $T_c$ ranges from 340°C to 480°C, consistent with magnetite. $T_c$ estimates in that samples with lower average measured Fe content have lower $T_c$ estimates. Of the T-type samples for which data are available, the $T_c$ estimates provide evidence for a population of oxides more Fe-rich than those measured via EMP. One (72.4°C h⁻¹) has two distinct Curie temperatures (240°C and 443°C), in spite of the fact that the EMP analyses cluster near the Mg, Al, and Cr end-members on the ternary plot (Figure 3). We interpret these to represent a chrome-rich (240°C) and a chrome-poor (443°C) phase, similar to those observed in natural, terrestrial chromite-bearing rocks [Radhakrishna Murthy and Krishnamacharyulu, 1994; Kumar and Bhalla, 1984]. The slowly cooled (5.7°C h⁻¹) TES-type sample has a Curie temperature (484°C) more consistent with the most magnetite-like of the EMP analyses, suggesting that the bulk of the unmeasured (smaller) crystals are more Fe-rich than the larger, measured crystals.

3.2.2. Hysteresis Data

Room temperature hysteresis measurements allow us to place constraints on the average magnetic grain size and domain state of the samples, including grains smaller than the resolution of the SEM. Within a given composition, $fO_2$ exerts the dominant control on grain size, while cooling rate plays a secondary role (Figure 4). Samples cooled under more oxidizing conditions and at slower cooling rates display more multidomain-like behavior. We note that this holds true even for T-type samples, which showed no obvious trends in oxide size with cooling rate in the EBS.
imagery. Some of the more reduced samples that plot along the SD-MD mixing trend show no visible oxides in the electron backscatter images or only show tiny, submicron oxides. These data are plotted as open symbols in Figure 4, and we suggest they represent mixtures of SD and superparamagnetic (SP) grain size populations, rather than SD-MD. Low-temperature remanence (e.g., Figures 5g and 5h and section 3.2.3) and susceptibility (Figures 6e and 6f and section 3.2.4) data are also indicative of a significant SP population.

As mentioned above, oxides in T-type samples visually appear smaller on average than in the M-type samples, and this distinction can also be seen in the hysteresis data. Interestingly, the data for the Cr-rich M-type samples are consistent with a much larger average magnetic grain size than the Cr-free samples of Brachfeld and Hammer [2006] (also shown in Figure 4). EBS imagery reveals that the maximum average grain size is indeed considerably larger in the Cr-rich samples. We attribute this to the presence of highly compatible Cr, which may increase the temperature at which an oxide phase is stable, thus enabling oxides growth to a larger size.

Hysteresis was also measured as a function of temperature for most of the QFM samples. More reduced samples were typically dominated by a paramagnetic signal, making calculation of hysteresis parameters difficult. All QFM samples, however, show a dramatic increase in coercivity below ~100 K that peaks at ~40–50 K for M-type samples and the most slowly cooled T-type sample, or continues to increase to the lowest measured temperature (10 K) for other T-type samples. This is likely related to a large increase in both magnetocrystalline anisotropy and magnetostriction in titanomagnetites at low temperatures [Schmidbauer and Readman, 1982; Moskowitz et al., 1998]. Saturation magnetization ($M_s$) decreases with decreasing temperature in all but one (variable cooling rate) of the meteorite-type samples; we interpret this as incomplete saturation of the titanomagnetite resulting in an overestima-

![Figure 4. Room temperature hysteresis data. Cooling rate shown in K h$^{-1}$ next to data points. Data from Cr-free samples from Brachfeld and Hammer [2006]. Open symbols likely contain a significant superparamagnetic contribution. Plot subdivisions for magnetite shown for reference.](image)

![Figure 5. SIRM imparted at 10 K or 20 K and measured on warming after cooling in 2.5 T field (FC, blue circles) and in zero field (ZFC, red circles). Note varying order of magnitude on magnetization scales. For both compositions, the FC remanence is always greater than (or equal to) the ZFC remanence, even in the most slowly cooled samples. This is in contrast to observations by Carter-Stiglitz et al. [2006], showing that the reverse is true for multidomain Fe$_{3-x}$Ti$_x$O$_4$ with x up to at least 0.35.](image)
tion of the paramagnetic component, and therefore an underestimate of $M_r$.

3.2.3. Low-Temperature Remanence (FC-ZFC)

[51] The low-temperature SIRM data for all QFM M-type samples are characterized by a large decrease in magnetization on warming at $\sim$50–60 K (Figure 5d), and the magnitude of the drop increases with decreasing cooling rate. The temperature of this transition ($T_{\text{crit}}$) is similar to that observed by Brachfeld and Hammer [2006] where it was interpreted as a magnetic isotropic point in low-Ti titanomagnetite [e.g., Moskowitz et al., 1998]. The cooling rate dependence may be related to grain size, and is similar to a grain size trend observed by Moskowitz et al. [1993] in synthetic SD magnetite. The pattern with respect to degree of cation substitution in the magnetite is consistent with that observed by Moskowitz et al. [1998]; samples determined to have more cation substitution on the basis of both EMP and Curie temperature analyses show little to no loss of remanence in the ZFC warming curve at temperatures below $T_{\text{crit}}$. Samples with less cation substitution show a more rapid decay in remanence on warming below $T_{\text{crit}}$. The T-type samples display a similar, though much less abrupt, remanence drop at slightly lower temperatures (Figure 5h). This pattern could imply a lower degree of cation substitution, which is consistent with the available $T_{\text{c}}$ data; however, the transition may be somewhat obscured by the gradual unblocking of a considerable SP component (also observed in the IW + 2 T-type samples, Figure 5g).

[52] In contrast to the QFM samples, all IW samples (except the T-type samples cooled at 231 and 72°C h$^{-1}$) have a phase that rapidly unblocks between 10 and 60 K (Figures 5a, 5b, and 5e). This is consistent with high-Ti hemoilmenite, which carries a high remanence in the spin glass phase at temperatures $\ll$50 K [Burton et al., 2008]. Additionally, all IW samples (except the M-type sample cooled at 18.7°C h$^{-1}$) have a phase that orders or blocks between $\sim$130 K and $\sim$150 K (Figures 5a, 5e, and 5f). The MELTS algorithm predicts both an ilmenite rich (90 mol %) rhombohedral phase and a Ti-rich spinel phase. The observed blocking at 130–150 K could be consistent with a spinel phase that is rich in Ti and other non-Fe cations.

[53] At least two other phases appear to be present in some of the IW samples. The M-type, 18.7°C h$^{-1}$ sample (as well as the M-type, IW +2 sample) has a phase that blocks at $\sim$220 K (Figures 5b and 5e). A peak in the susceptibility data (below and Figure 6a) at the same temperature is interpreted as a Néel point, indicative of an antiferromagnetic spin structure. This could be a spinel of a slightly different composition. We note that Cr-rich oxides are observed in the EDS data for this sample, and Fe-Cr spinels with ordering points in this temperature range would be antiferromagnetic [Robbins et al., 1971]. Again, it is likely that there are also other cations (Mg, Al) in the crystal structure, making a precise determination as to mineralogy difficult for all of these phases.

[54] Finally, the rapidly cooled (231 and 72°C h$^{-1}$) T-type samples (Figure 5f) have a phase that orders below 100 K. Interestingly, this phase appears to be antiferromagnetically coupled to the $\sim$150 K phase in the 72°C h$^{-1}$ sample (Figure 5f, solid circles), suggesting a two-phase intergrowth or another form of close coupling. Both of these samples (along with the M-type, IW, 231°C h$^{-1}$ sample) were suspended on Pt(-Fe) wire spirals during synthesis. Equivalent measurements carried out on a fragment of this wire from the 231°C h$^{-1}$ sample show some evidence for an ordering temperature at $\sim$100 K. This phase may therefore result from contamination of the sample with the wire or an iron oxide coating on the wire. A wire fragment from the 72°C h$^{-1}$ sample showed no evidence for ordering temperatures between 10 and 300 K, but did acquire a magnetization at 10 K.

Figure 6. (a–f) Frequency and temperature dependence of susceptibility. In-phase (out-of-phase) data shown in thick (thin) lines. Dashed line in Figure 6a is 4000 Hz data with estimated paramagnetic signal removed; scale on right-hand axis. M-type QFM samples at 231 and 19 K h$^{-1}$ (not shown) are similar to Figure 6c. T-type QFM samples at 231 and 5.7 K h$^{-1}$ are similar to Figures 6d and 6c, respectively. All IW samples not shown are similar to Figure 6d.
We note that the magnetic phases that are ordered below 130–150 K will not play a role in the magnetization of Mars’ crust. The minimum surface temperature at Mars’ poles is presently 150 K, warm enough that the inferred phases, if present on Mars, would behave paramagnetically. Only phases that are magnetically ordered with a nonzero moment (ferrimagnetic, canted antiferromagnetic) at and above surface temperatures are relevant to crustal magnetization.

3.2.4. Low-Temperature Susceptibility

M-type QFM samples display a frequency dependence in susceptibility between \( T_{c} \) (above) and \( 120–140 \) K, over which interval susceptibility increases sharply on warming (Figures 6b and 6e). The high-\( T_c \) (\( \sim 470\)°C) samples display a relatively large increase in susceptibility at or above \( T_{c} \) (Figure 6b), while the lower-\( T_c \) (\( \sim 330–340\)°C) samples show a smaller increase at \( T_{c} \) (Figure 6c); these patterns are similar to variations with composition observed in titanomagnetite [Moskowitz et al., 1998; Carter-Stiglitz et al., 2006]. Only the most slowly cooled T-type sample shows a hint of a similar transition. Data from more rapidly cooled samples (Figures 6e and 6f) are dominated by a broad peak showing frequency dependence characteristic of a large SP fraction [Worm, 1998; Worm and Jackson, 1999], or by a paramagnetic signal (231°C h\(^{-1}\), not shown).

All IW samples are characterized only by a paramagnetic signal (e.g., Figure 6d), with two exceptions. The M-type, 18.7°C h\(^{-1}\) sample (Figure 6a) has a peak at \( \sim 220 \) K, as mentioned in section 3.2.3. After removing the paramagnetic signal (by assuming it is similar to those in other M-type IW samples) the resulting curve (Figure 6a, dashed line) is reminiscent of an antiferromagnetic Néel point. Because this phase also carries a remanence (section 3.2.3), we assume it is canted or a defect moment.

The second exception is the T-type 231°C h\(^{-1}\) sample which shows evidence for a blocking temperature at \( \sim 200 \) K. Similar susceptibility behavior is observed a fragment of the Pt-Fe wire used for this sample, and as in the remanence-bearing phase that orders at \( 100 \) K (section 3.2.3), we interpret this as contamination from the wire.

3.2.5. Thermal Remanence

The intensity of TRM acquired is most strongly influenced by \( f_{O_2} \). In M-type samples, there is roughly linear relationship between \( \log_{10} f_{O_2} \) and \( \log_{10} \) TRM acquired in a 49 µT field (Figure 7a). We note that this trend likely cannot be extrapolated to higher \( f_{O_2} \), as larger, MD grains are not as efficient in acquiring a remanence; this can be seen in the work by Brachfeld and Hammer [2006], where TRM shows no dependence on \( f_{O_2} \) for \( f_{O_2} \geq \) QFM. The relationship is somewhat more complicated in T-type samples (Figure 7b). At \( f_{O_2} \leq 1 \)W+2 intensity is uniformly low (but above the noise level of the magnetometer). At QFM, remanence acquired varies over nearly 3 orders of magnitude, but the TRM of the most strongly magnetized samples is of the same order of magnitude as the M-type samples. Samples of both compositions that display no evidence for alteration with repeated heating show a linear relationship between applied field and acquired magnetization.

Median destructive temperatures (MDT) derived from thermal demagnetization of the TRM (Figures 7c and 7d) may provide additional constraints on oxide composition, although there is some evidence for sample alteration during the demagnetization. Weak-field susceptibility showed no consistent variation as the demagnetization progressed, but the duplicate TRM imparted after the complete demagnetization was considerably different for many samples. In general, QFM samples tended to show a slight reduction in acquired remanence, while more reduced samples showed an increase in remanence. The absolute magnitude of the change was a strong function of oxygen fugacity, with more reduced samples showing a greater degree of alteration, as might be expected since they are further from their “equilibrium” \( f_{O_2} \). We assume, however, that the alteration was progressive, and that the initial TRM acquired is the least affected. We denote samples with a >20% change in remanence with open symbols in Figures 7c and 7d. While the MDT of the more reduced samples especially must be treated with some skepticism, there are intriguing differences between the two bulk compositions. On average, with a decrease in \( f_{O_2} \), the MDT of the M-type samples decreases, suggesting an increase in non-Fe cations, while that of the T-type samples increases, suggesting a phase closer to that of pure magnetite.

3.2.6. Anhysteretic Remanence

ARM acquired at 100 mT AF in a 200 µT bias field is of the same order of magnitude as the TRM acquired in a 49 µT field, with no consistent offset in the positive or negative direction. The same trends in intensity with \( f_{O_2} \) are also observed. ARM acquisition is linear with applied DC field. Trends in cooling rate and composition can be seen in the coercivity spectra of the samples as determined by stepwise AF demagnetization of the ARM. Median destructive fields (MDF (Table 2)) of the T-type samples are consistently higher than those of the M-type samples, consistent with the hysteresis data and the observed smaller grain size. Additionally, within each composition at QFM, there is a trend of generally increasing MDF with increasing cooling rate, again corresponding to smaller average magnetic grain size at faster cooling rates. A strong outlier is the T-type sample cooled at 72.4°C h\(^{-1}\). Although one split of this sample is in line with the rest of the samples, a second split is considerably harder than all other measured samples, with an MDF of \( \sim 58 \) mT. The sample retains 22% of its magnetization at 100 mT AF. We note that the (few) EMP data available for this sample have the highest Cr content of any samples, and nonquantitative EDS data show that all analyzed grains contain significant Cr. Elevated coercivity has been linked to a Cr-rich spinel in natural rocks [Kumar and Bhatia, 1984; Radhakrishna Murthy and Krishnamacharyulu, 1994], and that may be the case here. MDF is more difficult to calculate accurately at lower \( f_{O_2} \) where the signal-to-noise ratio is considerably lower, and there are no obvious trends with cooling rate.

4. Discussion

The overall goal of this study is to understand how variations in bulk melt composition, \( f_{O_2} \), and cooling rate affect the resulting magnetic assemblage in a relatively rapidly cooled basalt. How do these factors control both the intensity of magnetization, as well as its stability with respect to temperature and time? Although the remanence carrying capacity of the samples is easy to measure, stability
with respect to temperature is largely related to composition, and stability through time is predominantly a function of oxide grain size. We will therefore briefly summarize our findings with respect to these two parameters before evaluating the role that our experimental variables play in controlling magnetic properties. We will also address how the resulting magnetic properties can be interpreted in terms of a potential contribution to the Mars magnetic source layer.

4.1. Oxide Composition and Implications for Crystal Nucleation and Growth

M-type QFM samples have analytical (EMP) data broadly consistent with the magnetism-based characterization of composition. The most slowly cooled samples and variably cooled M-type samples vary little in EMP analyses (Figure 3), having typical compositions of $\text{Fe}_{2.47}\text{Ti}_{0.17}\text{Al}_{0.17}\text{Mg}_{0.14}\text{Cr}_{0.02}\text{Mn}_{0.03}\text{O}_4$ and $\text{Fe}_{2.53}\text{Ti}_{0.12}\text{Al}_{0.15}\text{Mg}_{0.15}\text{Cr}_{0.01}\text{Mn}_{0.03}\text{O}_4$, respectively. The Curie temperature ($T_c$) of pure titanomagnetite with 2.5 apfu iron (i.e., $\text{Fe}_3\text{O}_4$) spinel ($\text{MgAl}_2\text{O}_4$) solid solution has $T_c \approx \text{280°C}$ [Dunlop and Özdemir, 1997]. However, the presence of other non-Fe cations can dramatically affect $T_c$. For example, an oxide with 2.5 apfu Fe in the magnetite ($\text{Fe}_2\text{O}_4$) spinel ($\text{MgAl}_2\text{O}_4$) solid solution has $T_c \approx \text{480°C}$ [Harrison and Putnis, 1996], which is very close to our measured $T_c$ of 468°C and 477°C for these two samples. In general we find that our measured $T_c$ and Fe values are consistent with data from the magnetite-spinel system; measured $T_c$ (Table 2) in samples cooled between...
Table 2. Sample Experimental Conditions and Magnetic Properties

<table>
<thead>
<tr>
<th>Sample</th>
<th>Desired ( f_O^2 )</th>
<th>Cooling Rate (deg C h(^{-1}))</th>
<th>Container Material</th>
<th>( f_O^2 ) Control(^a)</th>
<th>Batch(^b)</th>
<th>( T_B ) (deg C)</th>
<th>MDT(^c) (deg C)</th>
<th>MDF(^d) (mT)</th>
<th>( M_r ) (Am(^2) kg(^{-1}))</th>
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\(^{a}\) Method of controlling \( f_O^2 \) is flowing gases (G) or evacuated tube (T) with solid state buffer.

\(^{b}\) Oxide powder batch number.

\(^{c}\) Curie temperature.

\(^{d}\) Median destructive field.

\(^{e}\) Median destructive temperature.

\(^{f}\) Median destructive field.

\(^{g}\) Saturation remanent magnetization.

\(^{h}\) Exponentially varying cooling rate (see text).

\(^{i}\) Sample has anomalous texture, mineralogy, and magnetic properties suggesting experimental failure; magnetic properties are therefore not plotted in Figure 6.

\(^{j}\) Applying similar inferences to the T-type samples, Cr-rich grains and wide arrays of oxide compositions suggest rapid rates of oxide crystal growth at all cooling rates applied in this study. We speculate that the majority of oxide nucleation occurred at lower temperatures in the T-type runs, which is consistent with the presence of a magnetically ordered phase at room temperature.

231° C h\(^{-1}\) and 18°C h\(^{-1}\) are somewhat lower (340–400°C), consistent with a greater fraction of non-Fe cations. For many T-type QFM samples, we rely entirely on the magnetic data to infer the magnetically relevant oxide compositions. Curie and blocking temperature data suggest the presence of a population of (room temperature) remanence-carrying oxides that is on average more magnetite-rich than suggested by the EMP data. All samples have a phase that unblocks at temperatures >300°C (Figure 7), and the available (maximum) Curie temperatures are >440°C. Because the EMP analysis was conducted on only the largest crystals, we interpret EMP data as representing the high-temperature, early nucleated phases, with a higher proportion of non-Fe cations. Many of these larger crystals have EMP compositions inconsistent with a magnetically ordered phase at room temperature. Two samples additionally have a phase with low (<300°C) Curie or blocking temperatures (72°C h\(^{-1}\) and 19°C h\(^{-1}\)). These samples indeed have EMP compositions with high concentrations of Cr, but again with compositions that would be nonmagnetic at room temperature.

Compositional information on the more reduced samples (<QFM) is limited almost exclusively to blocking temperature data. Although these are associated with some uncertainty (see above), the data suggest that T-type samples have oxides that are on average more Fe-rich than the M-type counterparts. Additionally, reduced samples of both compositions appear to have both an ilmenite-rich rhombohedral phase and possibly a Ti- and/or Cr-rich spinel phase, neither of which will contribute to the magnetic source layer at crustal temperatures.
also exists a population that is considerably smaller, spanning into the SP range for T-type samples.

[71] While hysteresis data provide an average grain size, the full grain size distribution likely spans both SD and MD for many samples. It is therefore difficult to estimate how a remanence would decay through time, but we assume that samples with larger average grain size and more MD-like hysteresis behavior will tend to lose more of their remanence through time. We note that magnetic viscosity is commonly observed to be enhanced in coarse-grained rocks relative to fine-grained rocks; however, the theory of viscous magnetization (or demagnetization) has yet to be fully developed for MD grains, and at least one study has found that basalts with MD-sized ~TM55 do not appear to acquire a viscous magnetization [McClelland et al., 2003].

4.3. Variable Control on Magnetization

[72] Keeping in mind the compositional and grain size issues discussed above, we evaluate the role each of our variables plays in controlling the resulting magnetization. Figure 8 provides a summary of variation in magnetic parameters with two of our major experimental variables: oxygen fugacity and major element composition (predominantly Fe/Al ratio).

4.3.1. Oxygen Fugacity

[73] Within a given composition, fO2 appears to play the dominant role in controlling both oxide abundance and grain size, as well as magnetization, as previously observed [Brachfeld and Hammer, 2006]. Average grain size increases with increasing fO2 over the entire fO2 range investigated in this study, and continuing at least to IW+8.54 [Brachfeld and Hammer, 2006]. There is also a strong correlation between fO2 and TRM (Figure 7), as well as ARM, Mf, and Ms, especially for the Fe-rich, M-type samples. As noted in section 3.2.5, this correlation disappears above QFM for M-type samples [Brachfeld and Hammer, 2006], so ~QFM appears to be an optimum in terms of maximizing intensity while retaining an SD to PSD grain size.

[74] The role of fO2 on oxide composition is more difficult to determine, but it seems that the two bulk compositions may have opposite trends with respect to fO2. The room temperature remanence-bearing phase in M-type samples may become more Fe-poor with decreasing fO2, while that of the T-type samples remains relatively Fe-rich.

4.3.2. Bulk Composition

[75] The Fe/Al ratio of the sample appears to play a strong role in oxide abundance and grain size. M-type samples with high Fe/Al ratio (18 wt% FeO*) have a higher vol.% oxides and larger grain size than the low-Fe T-type samples. The magnetization is correspondingly higher in Fe-rich samples, as observed by Gee and Kent [1998]. TRM for the M-type samples is typically 1–2 orders of magnitude higher than that of the T-type samples. However, at QFM the most strongly magnetic T-type sample (72°C h⁻¹) is roughly equivalent to that of the M-type samples. Furthermore, the inclusion of highly compatible Cr in the melt seems to allow the nucleation of oxides at higher temperatures, allowing them to grow larger than in the absence of Cr.

4.3.3. Cooling Rate

[76] The cooling rate controls the size of the magmatic grains, which in turn controls the magnetic properties. For example, the cooling rate is likely to influence the degree of silicate crystallization, which in turn affects the magnetic properties. A slower cooling rate is likely to result in larger grain sizes and stronger magnetic properties, while a faster cooling rate is likely to result in smaller grain sizes and weaker magnetic properties.

Figure 8. Summary grid showing variation in magnetic properties with major experimental variables. Magnetic properties likely to influence the magnetic anomalies are magnetization intensity, blocking temperature (Tb), and domain state (SD, single domain; PSD, pseudo single domain; MD, multidomain). Likelihood that the given magnetic assemblage could result in an intense, stable magnetization, capable of producing the observed Martian magnetic assemblage is shown schematically by the shaded color bar: green is high probability, and white is low probability.

[68] The presence or absence of interstitial oxides is strongly linked to the degree of silicate crystallization in the T-type material (sections 3.1.2 and 3.1.3), indicating that the kinetics of silicate nucleation and growth largely dictate oxide abundance and composition. The T-type material has a much smaller solidification temperature interval, as assessed using the MELTS algorithm [Ghiorso and Sack, 1995]. Higher crystallinity of T-type samples for any given cooling trajectory (section 3.1.3) may reflect the influence of thermodynamic affinity on crystallization rate.

[69] The relative incoherence of mineralogies, phase proportions, and textures with respect to experimental variables in T-type experiments is consistent with the possibility that this composition is more sensitive to factors that influence nucleation kinetics, such as precise thermal history, starting material preparation, container roughness, etc. Plagioclase appears to be particularly prone to nucleation inhibition, a finding consistent with prior basalt dynamic cooling experiments [e.g., Grove, 1978; Grove and Bence, 1979], with the exception of observations by Lesher et al. [1999].

4.2. Oxide Grain Size

[70] Samples with a stable single domain grain size are more likely to have a magnetization that will persist through geologic time. For pure magnetite, the upper size limit for coherent, SD behavior is estimated at ~0.05–0.08 µm [Dunlop and Özdemir, 1997, and references therein]. For titanomagnetite of composition Fe₂₄Ti₀.₆O₄ (TM60), it is considerably larger: ~0.2 µm [Butler and Banerjee, 1975]. Both SD size estimates are far smaller than the largest grains observable in most of our samples, and indeed, the most slowly cooled samples have hysteresis dominated by MD or pseudo-single-domain (PSD) behavior. On the basis of room temperature hysteresis and low-temperature measurements, however, it is clear that for most other samples there...
The effect of cooling rate is most clearly seen within the QFM samples, where a clear trend in average magnetic grain size (or domain state) with cooling rate can be seen. This suggests that while a given $f_O^2$ may play a controlling role in the initial intensity of the magnetization, slower cooling rates may allow the grains to grow into an MD state, leading to a decrease in magnetization over geologic time scales.

Properties of the samples cooled with an exponential decreasing (or variable) cooling rate are inconsistent between the two bulk compositions. The M-type sample most closely resembles the more slowly cooled samples in oxide composition, abundance and magnetic properties. The T-type sample, by contrast, more closely resembles the more rapidly cooled samples in oxide abundance and magnetic grain size. (Compositions are harder to compare in the T-type samples.) Whether the magnetic properties are controlled by high-temperature rapid cooling or lower-temperature slower cooling is relevant as we consider cooling conditions of the bulk Martian crust. This variable cooling path was modeled to simulate the natural cooling path of a conductively cooled lava at 0.25 m from the surface; even if composed of thin flows, dikes or sills, the bulk of the crust almost certainly cooled more slowly than the model cooling path. A composition similar to our M-type composition (where magnetic properties are controlled by the low-temperature, slower cooling), might therefore be expected to develop even coarser oxides than those in any of our samples. A T-type composition, by contrast, may be able to retain a relatively small grain size even while cooling more slowly, because the magnetic properties would be most strongly influenced by the high-temperature, relatively rapid part of the cooling.

4.4. Modeling the Anomalies

To evaluate the possible contribution that each sample type might make to a magnetic anomaly, we construct a simple forward model. We magnetize a block that is square in cross section parallel to the planet’s surface and is 400 km on a side. The scale of the resulting anomalies is roughly that of some of the strong anomalies seen on Mars. Magnetization is calculated using the measured magnetization acquired in a 49 $\mu$T field (roughly Earth’s field) and assuming a density of 2900 kg m$^{-3}$. The magnetization is uniform in direction (vertical down). Intensity is allowed to vary with depth, following the blocking temperature distribution of the sample along a given (linear) geothermal gradient. This is done on a sample-by-sample basis by calculating temperature at a given depth, then finding the corresponding magnetization from the thermal demagnetization data shown in Figures 7e and 7f. We take the surface temperature of an early, warm Mars to be 273 K [Fanale et al., 1992]. The anomaly is calculated at 200 km altitude, using the method of a surface charge on a polygonal face [Blakely, 1996]. Results are shown (Figure 9) for samples produced under IW + 2 to QFM conditions, and for a range of geothermal gradients. We note that a reasonable gradient for early Mars may be $\sim$7–10 K km$^{-1}$ [Arkani-Hamed, 2000; Ruiz et al., 2006], although McGovern et al. [2002] suggest that it may have been $\sim$20 K km$^{-1}$ in places.

The solutions for most of the M-type QFM samples surpass the maximum observed anomalies on Mars for all reasonable geothermal gradients. More reduced samples, which have magnetizations equal to or nearly equal to the QFM samples, also have lower blocking temperatures; this results in thinner source layer and correspondingly decreased anomaly amplitude (dashed lines in Figure 9). Only one T-type sample results in modeled anomalies exceeding the necessary amplitude, and for smaller thermal gradients, which permit a thicker magnetized crust. The TES-type samples also tend to have a more distributed blocking temperature spectrum, which reduces the overall depth-integrated magnetization compared to the meteorite-type samples. We may conclude from this exercise that the more iron-rich meteorite sample can more easily acquire a magnetization both intense enough and with high enough blocking temperatures to produce the strong Martian anomalies, although we cannot entirely rule out a more Earth-like composition.

We note that these modeling results do not include the effects of magnetization in a smaller field, nonvertical magnetization, or viscous decay, all of which will reduce the anomaly amplitude. The effects of a smaller magnetizing field should be linear with applied field. A horizontal magnetization will decrease the anomaly amplitude by approximately one half, and Hood et al. [2007] found that isolated Mars anomalies are consistent with forward models magnetized predominantly in a north-south horizontal plane. The effects of viscous decay are harder to estimate, as discussed above, but if a significant fraction of the remanence is carried by multidomain grains, the decay will likely be considerable.

There are clear tradeoffs in achieving maximum intensity while retaining minimum grain size. Oxygen fugacity must be close to QFM in order to produce enough oxides, but in the M-type samples these oxides rapidly approach MD size as cooling rates approach geologically reasonable values. In the T-type samples, the average grain size remains smaller, but the most intensely magnetized samples have lower than average blocking temperatures, leading to a (relatively) strong magnetization, but a thinner total source layer. Restricting consideration to rapidly cooled materials, the most likely set of conditions to result in an intense, stable remanence may be an Fe-rich crust crystallized under moderately oxidizing (IW + 2.75 to QFM) conditions.

5. Conclusions

This paper reports on the results of an experimental examination of the magnetic assemblage produced under relatively rapid cooling in basalts of plausible Martian composition. We explore in detail the $f_O^2$ range most commonly found for basaltic meteorites as well as that over which a dramatic change in both oxide abundance and magnetic properties occurs. We systematically evaluate the effects of major element composition, oxygen fugacity ($f_O^2$), and cooling rate on phase chemistry and magnetic mineralogy, grain size (and domain state), and intensity of remanent magnetization. Figure 8 provides a summary of variation in magnetic parameters with two of our major experimental variables: oxygen fugacity and major element composition (predominantly Fe content).

Our results and conclusions are as follows.
The magnetic carrying phase in moderately oxidized (QFM) samples is a spinel-structured oxide with varying amounts of Cr, Ti, Mg, and Al. Under more reduced conditions (IW to IW+2) a room temperature remanence-bearing phase may be Fe-rich in TES-type samples and Fe-poor in meteorite-type samples.

Oxide compositions and compositional zonation are strongly dependent upon cooling rate in Fe-rich magma, but less so in Al-rich magma, because of differences in crystallization sequence and the kinetics of silicate mineral nucleation and growth.

Oxygen fugacity exerts the primary control on oxide grain size and abundance. Additionally, an Fe-rich melt leads to more, larger oxides, and the presence of Cr appears to stabilize oxides at higher temperatures, leading to larger oxides. Slower cooling rates also lead to larger average grain size.

Moderately oxidized (~QFM) samples of both compositions can acquire an intense thermal remanent magnetization.

However, a volcanically generated Mars crust of Fe-rich (meteorite-type) composition can more easily produce large magnetic anomalies over a wide range of geothermal gradients.

Increase of grain size into the multidomain range or low blocking temperatures may decrease anomaly amplitude and/or affect long-term stability of remanence.

In rapidly cooled materials, the most likely set of conditions to result in an intense, stable remanence may be an Fe-rich crust crystallized under moderately oxidizing (IW + 2 to QFM) conditions.

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